# Causal Structure of Multiboundary Wormholes

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## Introduction

Multiboundary wormholes in 2+1 dimensions

- Natural generalization of eternal BTZ black hole
- Quotients of  $AdS_3 \rightarrow \text{locally } AdS_3$
- Traversable multiboundary wormholes?

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# Eternal BTZ black hole

[Maldacena '01] Eternal AdS black holes ↔ Thermofield double state (TFD)

$$H = H_R - H_L$$
  

$$Z = \operatorname{Tr} e^{-\beta H}$$
  

$$CFT_L$$
  

$$t_L = -t$$
  

$$CFT_R$$
  

$$t_R = t$$

Thermofield double state

$$|TFD\rangle = \frac{1}{\sqrt{Z}} \underbrace{\sum_{n} e^{-\beta E_{n}/2} |E_{n}\rangle_{L} |E_{n}\rangle_{R}}_{\text{entangled, non-interacting}}$$

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Traversable wormhole via double trace deformation

Double trace interaction [Gao, Jafferis, Wall '16]



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# Multiboundary wormholes

• Einstein gravity in 2+1 dimensions with negative cosmological constant:  $R_{\mu\nu\rho\sigma} = \Lambda(g_{\mu\rho}g_{\nu\sigma} - g_{\mu\sigma}g_{\nu\rho})$ 

 $\Rightarrow$  Spacetime geometry is locally  $AdS_3$ 

• BTZ is quotient of  $AdS_3$ .



Multiboundary wormholes are generalizations of eternal BTZ.

 $\mathcal{M} = AdS_3/\Gamma$ ,  $\Gamma$  discrete subgroup

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# Causal structure of BTZ as a quotient

**Time dependence**: Identified geodesics evolve into geodesic planes [Aminneborg et al '97, Brill '99]



Image: A matrix

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**Singularites**: Intersection of geodesic planes **Horizon**: Backwards light cone of point *P* 

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# Multiboundary wormhole

• Three-boundary wormhole (genus zero)



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- Curves of minimal length between the identified geodesics characterize the horizons
- **Causal shadow**: Region causally disconnected from all asymptotic boundaries.

Time development of three-boundary wormhole [Aminneborg et al '97, Brill '99]





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# Fenchel-Nielsen coordinates

Cut Riemann surface with into pairs of pants [Skenderis, Rees '09]

m = number of boundaries g = genus



- Assign lengths  $L_i$  to all edges of every pants
- Assign a twisting parameter  $t_j$  for every gluing of two pair of pants

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## Fenchel-Nielsen coordinates



- (a) Pick up a closed geodesic
- (b) Thickening geodesic to obtain cylinder
- (c) Extend cylinder as far as possible until bounding circles self intersect.

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#### Outer chart coordinates

- Outer chart: covers one asymptotic region and part of the interior
- Explicit metric [Skenderis, Rees '09]

$$ds^{2} = \frac{\rho^{2} + M}{\cosh^{2}(\sqrt{M}\tau)}(-d\tau^{2} + d\phi^{2}) + \frac{d\rho^{2}}{\rho^{2} + M}$$

- Horizons lie at the surfaces  $\rho = \sqrt{M} |\sinh(\sqrt{M}\tau)|$
- Exterior is exactly BTZ

$$ds_{\mathsf{BTZ}}^{2} = -(r^{2} - M)dt^{2} + \frac{dr^{2}}{r^{2} - M} + r^{2}d\phi^{2}$$

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# Outer chart coordinates

- Singularity r=0 corresponds to  $\tau \to \pm \infty$
- Boundaries at  $\rho = \pm \infty$

$$\frac{\cosh(\sqrt{M}\tau)\rho}{\sqrt{\rho^2 + M}} > -\sqrt{\frac{C'^2 + C''^2 + 2C'C''C}{C'^2 + C''^2 + C^2 + 2C'C''C - 1}}, \quad C' = \cosh(\pi\sqrt{M'})$$



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## Transition functions

$$\rho' = -\sqrt{\frac{M'}{M}} \left( \cosh(A)\rho + \sinh(A)\cosh\left(\sqrt{M}(\phi - f)\right) \frac{\sqrt{\rho^2 + M}}{\cosh(\sqrt{M}\tau)} \right)$$
$$\sqrt{M} \tanh(\sqrt{M'\tau'}) \sqrt{\rho'^2 + M'} = \sqrt{M'} \tanh(\sqrt{M}\tau) \sqrt{\rho^2 + M}$$
$$e^{2\sqrt{M'}(\phi' - f')} = \frac{\rho\cosh(\sqrt{M}\tau) + \sqrt{\rho^2 + M}\cosh(\sqrt{M}(\phi - f) - b)}{\rho\cosh(\sqrt{M}\tau) + \sqrt{\rho^2 + M}\cosh(\sqrt{M}(\phi - f) + b)}$$

Transition from  $1 \rightarrow 2$  chart (unprimed to primed) has  $f = 0, f' = \pi$ .

$$\cosh(A) = \frac{\cosh(\pi\sqrt{M})\cosh(\pi\sqrt{M'}) + \cosh(\pi\sqrt{M''})}{\sinh(\pi\sqrt{M})\sinh(\pi\sqrt{M'})}$$
$$\sinh(A)\sinh(b) = 1$$

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Null geodesics in three-boundary wormhole

Null curve starting from past horizon of chart  $(U_1, V_1)$ 



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- Shift in  $U_2$  direction in chart  $(U_2, V_2)$
- Dependence on angular variable

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# To do

- More complete characterization of causal shadow
- Bulk-to-boundary propagators
- Construction of a traversable three-boundary wormhole
- Extension to more boundaries, non-zero genus

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